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REVISION OF GEODETIC PARAMETERS

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(NASA-CR-143331) REVISION OF GEODETIC
PARAMETERS (Smithsonian Astrophysical
Observatory) 26 p HC \$3.75 CSCL 08N

N75-30696

Unclas
G3/46 33052

August 1975



Smithsonian Institution
Astrophysical Observatory
Cambridge, Massachusetts 02138



NG 209-015-002

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ABSTRACT

Laser data from nine satellites and 12 stations are combined with surface-gravity data to obtain spherical harmonics representing the geopotential complete through degree and order 18. This laser-data-only solution provides a reasonable improvement to the gravity field.

INTRODUCTION

Smithsonian Astrophysical Observatory has published gravity-field solutions utilizing both satellite-tracking and terrestrial gravity data (see, e.g., Gaposchkin, 1974), which were based primarily on precision-reduced camera data and on then-available surface-gravity data. Now, however, we have better data, and new types of data will soon be available. The work reported here is the beginning of a general revision and extension of our knowledge of the geopotential. In this first iteration, we will experiment with new data (laser ranges), verify new methods of data reduction, and prepare for new types of data (altimeter and satellite-to-satellite tracking data). We seek here to improve our knowledge of the gravity field so that better satellite orbits, consistent with surface-gravity data, can be calculated. This is done by using laser data only, whose accuracy (approximately 1 m) is far greater than the ephemeris accuracy (approximately 10 m). The situation has changed since 1971, when the bulk of the data used were 4-arcsec camera data. Furthermore, substantial revisions have been made in the treatment of orbit perturbations, and all these

advances indicate a new attempt at improved geopotential coefficients.

OBJECTIVES

This computation has several objectives, the primary one being to use laser data only in a determination of the earth's gravity field, with the aim of computing satellite orbits to an accuracy comparable with that of the laser data. We can obtain a realistic gravity field consistent with surface-gravity data. A secondary objective is to study the consequences of using data that provide no ties to an inertial reference frame, as was the case with camera data. For our third objective, we will investigate the effects of one satellite on another in the solution; that is, we can optimize a solution for one satellite by using only data from that satellite. The question then becomes: How much does adding data from a second satellite degrade the orbit computed for the first? Of course, improved orbits and a more accurate geoid are necessary for analyzing satellite altimetry data for geodetic and oceanographic purposes.

REFINEMENTS AND TECHNIQUES

Improvements in perturbation calculations have been numerous. The inclination function for tesseral harmonics, as formulated by Kaula, computationally loses accuracy for high-degree coefficients. It has been replaced by a mathematically equivalent formula derived from group theory (Gaposchkin, 1973). The interaction terms between J_2 and resonant harmonics have also been improved. Lunar and solar perturbations and body tides and ocean tides have been computed to the necessary accuracy (Kozai, 1973). Perturbations arising from the noninertialness of the coordinate system have been corrected and improved (Kinoshita, 1975a,b),

and those due to direct solar radiation pressure (Aksnes, 1975), albedo pressure (Lautman, 1975a,b), and infrared radiation (Lautman, 1975c) have all been included and tested (Gaposchkin et al., 1975).

The compilation of surface-gravity data used in Gaposchkin (1974) has since been augmented (Williamson and Gaposchkin, 1975). The surface-gravity data are summarized in Table 1, and the distribution of these data is shown in Figure 1. Table 2 compares the $1^\circ \times 1^\circ$ gravity anomalies from the Defense Mapping Agency Aerospace Center (DMAAC, 1973) with other available compilations.

Three coordinated programs have provided laser tracking data in sufficient density to be used for a gravity-field determination: the International Satellite Geodesy Experiment (ISAGEX) held in 1971, the Earth Physics Satellite Observation Campaign (EPSOC) in 1972-1973, and the Geos 3 campaign in 1975. There has been a steady improvement in the volume, reliability, and accuracy of the data. One of the objectives of the ISAGEX program was to obtain data for determining the gravity field, and a number of orbital arcs are suitable for this purpose. The EPSOC program was directed to the study of long-period effects and polar motion. Some arcs from EPSOC are also used for determining the tesseral harmonics. Finally, during the current Geos 3 program routine data are being obtained and included in the analysis. Table 3 lists the participating stations for each campaign that are used in the analysis reported here. The distribution of the stations is shown in Figure 2.

Currently, nine satellites in orbit are equipped with cube-corner reflectors suitable for laser ranging; these satellites are useful for a gravity-field determination, especially as their distribution in inclination and height is reasonably good. Table 4 lists their orbital characteristics and the number of arcs used, and the distribution of satellites is given in Figure 3.

Satellite orbits are usually computed for between 8 and 16 days, the interval depending primarily on the availability of data. Also, each arc covers at least one oscillation of the resonant period. Table 5 gives the resonant periods for the satellites, while Table 6 lists the constants used to calculate the satellite orbits.

SOLUTIONS

The normal equations for surface-gravity data have been computed complete from degree 2 through 18. The combination solution included a number of harmonics of higher degree that are resonant with one or more of the satellite orbits. To this set of surface-gravity normal equations was added a set of normal equations, satellite by satellite. The system was solved after adding one satellite (Geos 1), and the result has been compared with all the satellites and with surface gravity. For the satellite comparison, one arc from each satellite was selected: generally, the arc with the most data. The orbit was recalculated with the revised gravity field, and the orbital fit in terms of σ_0 was used as the criterion. In all combinations, the satellite data were used at their a priori weight. The surface gravity was given several weights, and all solutions were tested in order to determine the optimum weight for the combination.

The weight finally adopted is

$$< A > \frac{27}{nA} \text{ mgal} ,$$

where n is the number of $1^\circ \times 1^\circ$ squares in each $5^\circ \times 5^\circ$ mean, A is the area of the anomaly, and $< A >$ is the average area. This weight is twice that used in the 1973 Smithsonian Standard Earth (III) (SE III) (Gaposchkin, 1973). The ISAGEX laser data were given a 5-m weight, and all other laser data, a 2-m weight.

These combination solutions are summarized in Tab' 7. The orbit for a satellite not used in the solution is really very poor, generally because of relatively small changes in a few coefficients resonant with the gravity field. We note that

Peole, which has no resonance, degrades marginally. The surface gravity $\langle(\Delta g_t - \Delta g_s)^2\rangle$ is only relative. A large part of those residuals is information in the higher harmonics. The best fit is obtained with only one satellite; however, when satellite data are added, the degradation is not large considering the overall accuracy of surface-gravity data.

Generally, adding satellite data also degrades the satellite orbit fit, but not very much. This overall improvement is considered quite satisfactory for one iteration of a very complex nonlinear process. All satellite orbits improved by at least 1 m^2 in the orbital fit. In percentage terms, Geos 1, BE-C, Geos 2, Starlette, Peole, and Geos 3 each improved in orbital fit by 17.9, 22.2, 43.9, 69.2, 26.2, and 13.6%, respectively, for an average 32% improvement!

The final adopted 52-arc solution (SE IV.1) can be compared with the surface-gravity data in more detail. Assuming they are statistically independent, the following quantities defined by Kaula (1966) can be computed and used to compare a geopotential model (g_s) with observed values of surface gravity (g_t):

$\langle g_t^2 \rangle$	The mean value of g_t^2 , where g_t is the mean free-air gravity anomaly based on surface gravity, indicating the amount of information contained in the surface-gravity anomalies.
$\langle g_s^2 \rangle$	The mean value of g_s^2 , where g_s is the mean free-air gravity anomaly computed from the geopotential model, indicating the amount of information in the computed gravity anomalies.
$\langle g_t g_s \rangle$	An estimate of g_h - i.e., the true value of the contribution to the gravity anomaly of the geopotential model and the amount of information common to both g_t and g_s .
$\langle (g_t - g_s)^2 \rangle$	The mean-square difference of g_t and g_s .
$E(\epsilon_s^2)$	The mean-square error in the geopotential model.
$E(\epsilon_t^2)$	The mean-square error of the observed gravity.
$E(\delta g^2)$	The mean square of the error of omission - that is, the difference between true gravity and g_h ; this term is then the model error.

If the geopotential model were perfect, then $\langle g_s^2 \rangle = \langle g_h^2 \rangle$, which in turn would equal $\langle g_t g_s \rangle$ if g_t were free from error and known everywhere. Then, ϵ_s^2 would be zero even though g_s would not contain all the information necessary to describe the total field. The information not contained in the model field — i.e., the error of omission, δg — then consists of the higher order coefficients. The quantity $\langle (g_t - g_s)^2 \rangle$ is a measure of the agreement between the two estimates g_t and g_s and is equal to

$$\langle (g_t - g_s)^2 \rangle = E(\epsilon_s^2) + E(\epsilon_t^2) + E(\delta g^2) .$$

Another estimate of g_h can be obtained from the gravimetric estimates of degree variance σ_ℓ^2 (Kaula, 1966):

$$E(g_h^2) = D = \sum_{\ell} \frac{n_{\ell}}{2\ell + 1} \sigma_{\ell}^2 ,$$

where n_{ℓ} is the number of coefficients of degree ℓ included in g_h , and

$$\sigma_{\ell}^2 = \gamma^2 (\ell - 1)^2 \sum_m (\bar{C}_{\ell m}^2 + \bar{S}_{\ell m}^2) .$$

We also have

$$E(\epsilon_s^2) = \langle g_s^2 \rangle - \langle g_s g_t \rangle$$

and

$$E(\epsilon_t^2) = \langle g_t^2 \rangle / \langle n \rangle .$$

These values are given in Table 8 for SE III and for this solution.

The information in the surface-gravity data solution $\langle g_t^2 \rangle$ has increased in this new data set. This is reasonable; since the unobserved areas had an expected value of zero, the fewer observed areas there are, the lower the variance is. However, the information in the satellite solution $\langle g_s^2 \rangle$ has decreased, a fact that is confirmed by a decrease in D. Therefore, the information in SE III was too high. The residual $\langle (g_t - g_s)^2 \rangle$ has remained roughly the same, while the information in the higher harmonics is estimated to be larger. The estimate of $E(\epsilon_s^2)$ cannot be good, as these sets of data, g_s and g_t , are not independent.

The spherical-harmonic coefficients are listed in Tables 9 and 10. Figure 4 is a plot of the mean potential coefficient as a function of degree, and Figures 5 and 6 show the geoid height and gravity anomalies for this solution.

FUTURE WORK

The obvious next step is to complete another iteration, taking the solution to perhaps degree and order 24. To this can be added the normal equations for zonal harmonics and sets of resonant harmonics. When the orbital accuracy approaches a few meters, then we must reduce the error in the station coordinates by solving for them. Finally, of course, the aim is to add altimetry data to this system of normal equations.

In summary, we find that the improved accuracy that has been apparent from laser data is becoming realized.

ACKNOWLEDGMENT

This work was supported in part by grant NGR 09-015-002 from the National Aeronautics and Space Administration.

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Table 1. Surface-gravity data available.

Source	No. of 1° x 1° means	No. of 550 x 550 km blocks	
		n ≥ 1	n = 25
Gaposchkin (1974)	19328	1183	145
Williamson and Gaposchkin (1975)	31636	1452	485
Maximum number	64800	1690	1690

Table 2. Comparison of $1^{\circ} \times 1^{\circ}$ mean gravity anomalies with DMAAC (1973).

Source	No. of points compared	Mean difference (mgal)	rms (mgal)
Australia (Mather, 1970)	1364	1.64	24.16
North America and North Atlantic (Talwani <i>et al.</i> , 1972)	3613	-0.18	15.29
Indian Ocean (Kahle and Talwani, 1973)	2226	-1.66	23.09
Worldwide (ACIC, 1971)	19164	-0.23	16.99

Table 3. Lasers stations used in this analysis.

Station			ISAGEX	EPSOC	Geos 3
Number	Location	Agency			
7902	Olifantsfontein, S. Africa	SAO	x	x	x
7907	Arequipa, Peru	SAO	x	x	x
7921	Mt. Hopkins, Arizona	SAO	x	x	x
7928	Natal, Brazil	SAO	x	x	x
7930	Athens, Greece	SAO	x		x
7050	GSFC	NASA	x		x
7060	Guam	NASA	x		
7061	San Diego, Calif.	NASA		x	
7080	Quincy, Calif.	NASA		x	
7068	Grand Turk Island	NASA			x
7804	San Fernando, Spain	CNES	x		
7809	Haute Provence, France	CNES	x		

Table 4. Summary of dynamical data.

Satellite		Inclination	Eccentricity	Perigee height (km)	a (km)	Number of arcs
Designation	Name					
7010901	Peole	15°	0.017	635	7070	5
6701401	D1D	39	0.053	569	7337	3
6701101	D1C	40	0.052	579	7336	2
6503201	BE-C	41	0.026	941	7311	9
7501001	Starlette	50	0.0207	805	7335	5
6508901	Geos 1	59	0.073	1121	8074	14
7502701	Geos 2	-65	0.0005	840	7222	4
6800201	Geos 3	-75	0.031	1101	7709	8
6406401	BE-B	80	0.012	912	7362	2
Total						52

Table 5. Resonant periods.

Satellite	Resonant with order m	Period (days)
7010901	none	
6701401	13	9.4
6701101	14	2.6
6503201	13	5.6
7501001	14	3.2
6508901	12	7.2
7502701	14	3.9
6800201	13	6.3
6406401	14	2.9

Table 6. Constants used in orbit computation.

$GM = 3.986013 \times 10^{20} \text{ cm}^3 \text{ sec}^{-2}$
$c = 2.997925 \times 10^{10} \text{ cm sec}^{-1} = \text{speed of light}$
$k_2 = 0.25 = \text{Love's number}$
$\epsilon_2 = 10^\circ = \text{phase lag of tide}$
$a_e = 6.378140 \text{ Mm}$
$\alpha = 0.32 = \text{earth's albedo}$

Satellite	A/m ($\text{cm}^2 \text{ g}^{-1}$)
7010901	0.20
6701401	0.30
6701101	0.30
6503201	0.13
7501001	0.01
6508901	0.10
7502701	0.04
6800201	0.06
6406401	0.10

Table 7. Comparison of solutions.

Solution	No. of arcs	Surface gravity $\langle (\Delta g_t - \Delta g_s)^2 \rangle, \ell \leq 18$ (mgal ²)		<u>Geos 1</u>		<u>BE-C</u>		<u>Geos 2</u>		<u>Starlette</u>		<u>Peole</u>		<u>Geos 3</u>	
		$n \geq 1$	$n \geq 20$	σ_0	n	σ_0	n	σ_0	n	σ_0	n	σ_0	n	σ_0	n
SE III	(203)	193	135	4.53	3021	5.81	1699	5.58	1112	16.61	2443	15.41	805	13.16	1076
Geos 1	(5)	123	86	3.15	3033	38.87	1688	32.60	1126	24.14	2336	16.16	797		
and BE-C	(5)	130	89	3.12	3004	4.22	1701	133	1140	770	2344	12.35	796		
and Geos 2	(5)	132	91	3.38	3012	4.20	1689	3.23	1125	18.11	2316	12.59	794		
Geos 1 BE-C	(14) (9)														
Geos 2	(8)														
BE-B	(2)	139	99	3.58	3011	4.50	1696	3.32	1130	4.96	2259	12.40	807		
D1D	(3)														
D1C	(2)														
Peole	(5)														
Starlette	(5)														
and Geos 3	(4)	142	100	3.72	3022	4.52	1695	3.13	1122	5.12	2258	11.38	794	11.37	1075

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Table 8. Comparison of surface gravity with solution (mgal²); all solutions are to the 18th degree and order.

Solution	$\langle (g_t - g_s)^2 \rangle$	$\langle g_t g_s \rangle$	$\langle g_s^2 \rangle$	D	$\langle g_t^2 \rangle$	$E(\epsilon_s^2)$	$E(\epsilon_t^2)$	$E(\delta_g^2)$	n*
SE III+	156	202	258	237	302	56	24	75	n \geq 1 (1183)
SE III+	117	255	281	237	345	26	19	63	n \geq 10 (659)
SE III+	105	221	236	237	311	15	13	77	n \geq 20 (301)
SE III	135	237	254	235	356	16	15	104	n \geq 20 (678)
SE IV.1	142	192	217	227	309	25	19	98	n \geq 1 (1452)
SE IV.1	103	232	235	227	333	2	16	85	n \geq 10 (1056)
SE IV.1	100	250	244	227	356	-6	15	92	n \geq 20 (678)

* n is the number of 1°x1° mean gravity anomalies in the 5°x5° mean gravity anomalies.

+ From the available data, there were 1183, 659, and 306 gravity anomalies with n = 1, 10, and 20 1°x1° anomalies.

Table 9. SE IV.1 fully normalized tesseral-harmonic coefficients for the geopotential (in units of 10^{-6}).

L	M	C	S	L	C	S	L	M	C	S
2	2	1.3390	1.3392	3	2	-26870	3	2	74764	72826
3	3	1.5333	-1.5333	4	1	-57429	4	2	39558	-65655
4	4	96673	17236	4	4	-26977	5	1	-56333	12083
5	2	63513	-37943	5	3	-53617	5	4	-29932	-33202
5	5	61770	61273	6	1	-67106	6	2	39469	54533
6	3	76855	-36232	6	4	-65148	6	5	-19537	-13138
6	6	63578	24634	7	1	-25179	7	2	22445	-05844
7	3	20820	18135	7	4	16993	7	5	01853	-05908
7	7	31653	-15875	8	3	06660	8	1	-06653	-15064
8	2	13553	-03632	8	6	-03379	8	4	07779	-13619
8	8	03193	03713	9	3	-05563	9	2	-03315	03947
9	5	-14459	-04102	9	4	13493	9	5	-10207	01352
9	9	-16767	04936	9	7	-02758	10	2	03992	00369
9	9	04255	-18792	10	1	-11769	10	8	01939	05077
10	3	-11362	11390	10	4	08325	10	5	-03992	00369
10	6	-11755	05235	10	7	-12391	10	8	00867	07231
10	10	-17334	10090	10	10	01450	11	1	-03992	-01015
11	2	11610	-00077	11	3	12283	11	4	00520	20579
11	5	-04242	09413	11	6	-04767	11	7	09059	10093
11	8	06461	-03083	11	9	03615	11	10	-04845	04960
11	11	-04100	06399	12	1	-00443	12	2	-00339	-00595
12	3	11844	-01694	12	4	-09670	12	5	09202	-03075
12	6	04415	-07292	12	7	-12473	12	8	-02826	-07762
12	9	-06147	05012	12	10	-04730	12	11	02373	00708
12	12	-00076	-07439	13	1	-06572	13	2	00804	04858
13	3	00085	01429	13	4	09041	13	5	04432	-04956
13	6	-02107	-05285	13	7	03022	13	8	00277	04741
13	9	-04216	00330	13	10	04597	13	11	-06226	-04713
13	12	-02344	-09545	13	13	-06281	14	1	-01211	-10001
14	2	-00741	01453	14	3	09587	14	4	00531	03879
14	5	-01203	02788	14	6	-01590	14	7	02242	04756
14	8	-04135	00275	14	9	03968	14	10	02462	00118
14	11	03814	04464	14	12	00863	14	13	04310	-04120
14	14	-06698	-02118	15	1	04914	15	2	-00750	02813
15	3	-00114	01633	15	4	-02792	15	5	02160	-02448
15	6	-02151	08104	15	7	08711	15	8	-05521	-03260
15	9	-03379	-01423	15	10	-04090	15	11	00194	03595
15	12	-01924	-04109	15	13	-02269	15	14	00367	03905
15	15	00176	-07630	16	1	-01224	16	2	00328	-06034
16	3	-01669	-02335	16	4	07270	16	5	-02712	-03154
16	6	03896	02730	16	7	-01526	16	8	-02338	-03223
16	9	03628	04728	16	10	00312	16	11	00552	-03810
16	12	02270	00997	16	13	01407	16	14	-01578	03977
16	15	-06156	12601	16	16	-00795	17	1	-01884	02527
17	3	-04415	-00175	17	3	03634	17	4	-05198	-04764
17	5	01310	01175	17	6	-06345	17	7	01898	01249
17	8	00076	02580	17	9	-03490	17	10	01920	-04102
17	11	-01010	00825	17	12	04423	17	13	01422	-03708
17	14	-06139	-00654	17	15	02338	17	16	-06713	-02418
17	17	00613	01562	18	1	-00981	18	2	01288	-00977
18	3	-05049	01711	18	4	00361	18	5	00067	02462
18	6	01314	-03672	18	7	02995	18	8	04397	00422
18	9	01015	-03672	18	10	-02995	18	11	-02204	-03591
18	12	-03691	01458	18	13	-00817	18	14	00711	05891
18	15	-03541	02154	18	16	04460	18	17	03483	01169
18	18	-01377	-01230	19	1	02980	19	13	01157	02736
19	14	00442	00763	20	13	05582	20	14	05653	-06967
21	13	01760	00554	21	14	02035	22	13	00068	-00427
22	14	00124	04693	23	13	00601	23	14	00716	00903
24	14	04139	-03852							

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Table 10. SE III fully normalized zonal harmonic coefficients for the geopotential (in units of 10^{-6}).

L	M	C	L	M	C	L	M	C	L	M	C	L	M	C
2	0	-484.16999	3	0	.96041	4	0	.53933	5	0	.06874			
6	0	-.15310	7	0	.09089	8	0	.04972	9	0	.03533			
10	0	-.05172	11	0	-.06506	12	0	.03840	13	0	.06524			
14	0	-.01950	15	0	-.01886	16	0	-.00592	17	0	.03719			
18	0	.01677	19	0	-.01585	20	0	.01858	21	0	.01266			
22	0	-.01371	23	0	-.02115	35	0	.01590	36	0	-.02329			

FIGURE CAPTIONS

- Figure 1. Distribution of $1^\circ \times 1^\circ$ mean surface-gravity data.
- Figure 2. Locations of the observing stations included in SE IV.1.
- Figure 3. Distribution of perigee heights and inclinations of the satellites used in SE IV.1.
- Figure 4. Mean potential coefficient by degree.
- Figure 5. SE IV.1 geoid height in meters calculated with respect to the best-fitting ellipsoid, $f = 1/298.256$.
- Figure 6. SE IV.1 gravity anomalies in milligals calculated with respect to the best-fitting ellipsoid, $f = 1/298.256$.

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Figure 1. Distribution of $1^{\circ} \times 1^{\circ}$ mean surface-gravity data.

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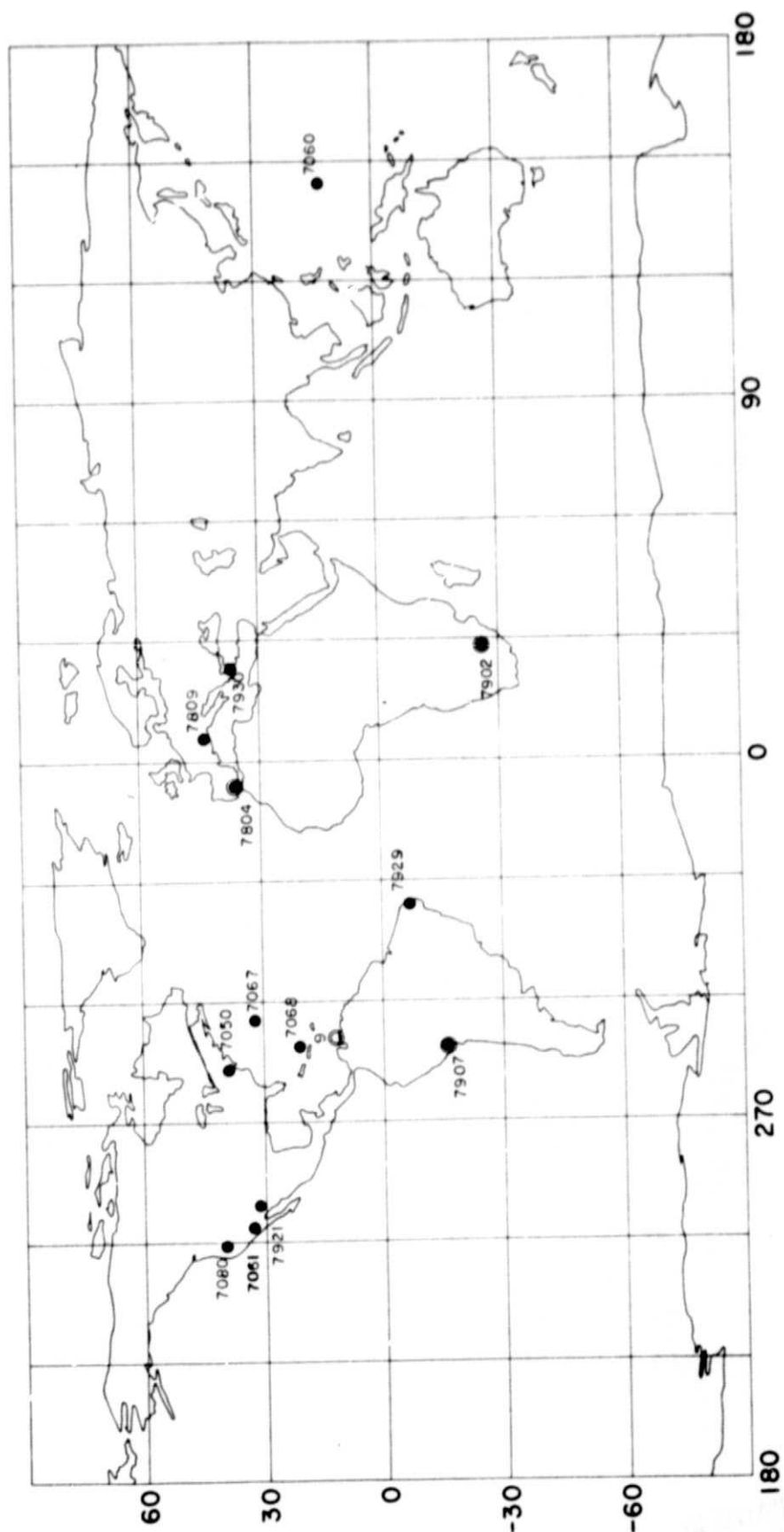


Figure 2. Locations of the observing stations included in SE IV.1.

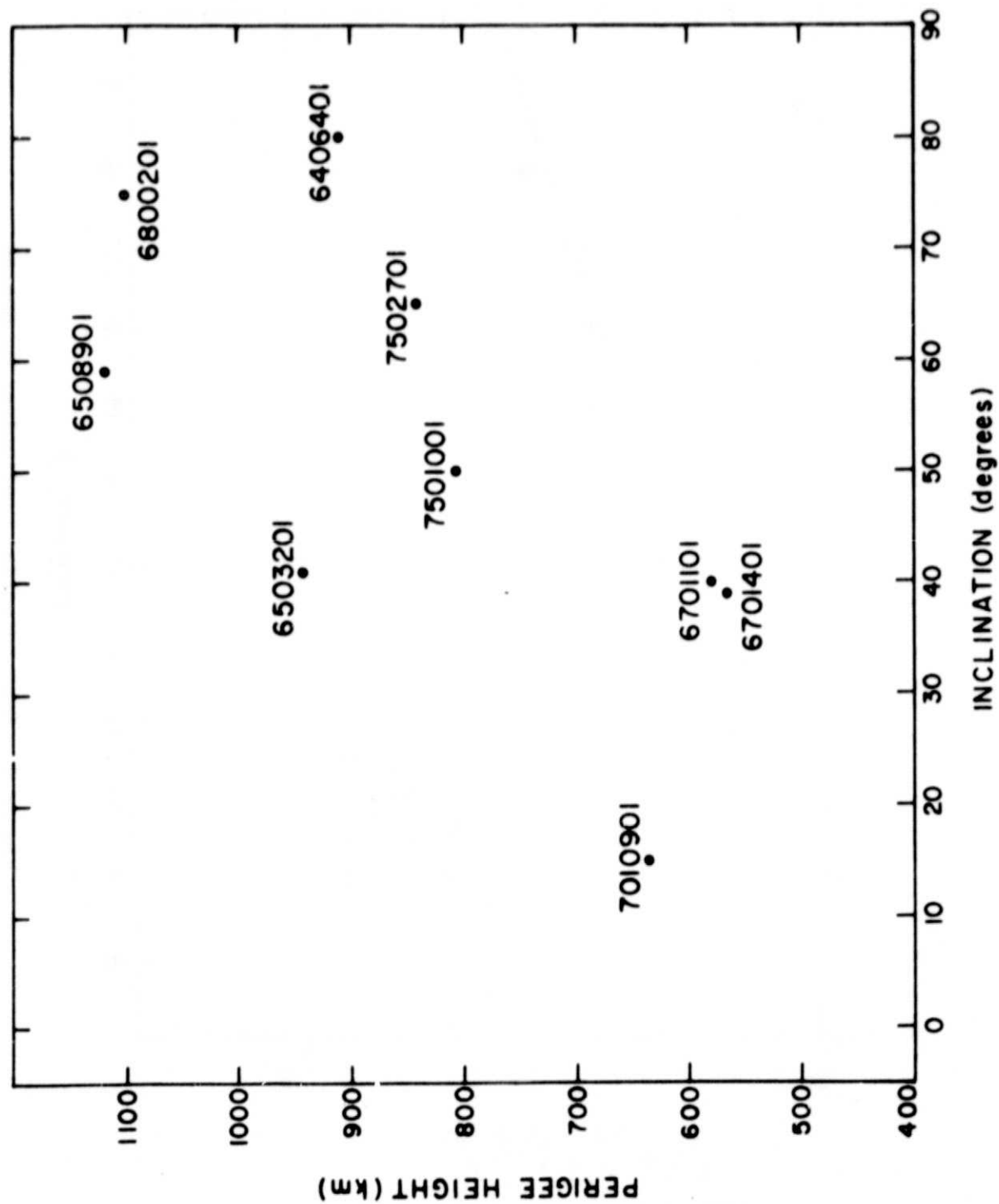


Figure 3. Distribution of perigee heights and inclinations of the satellites used in SE IV.1.

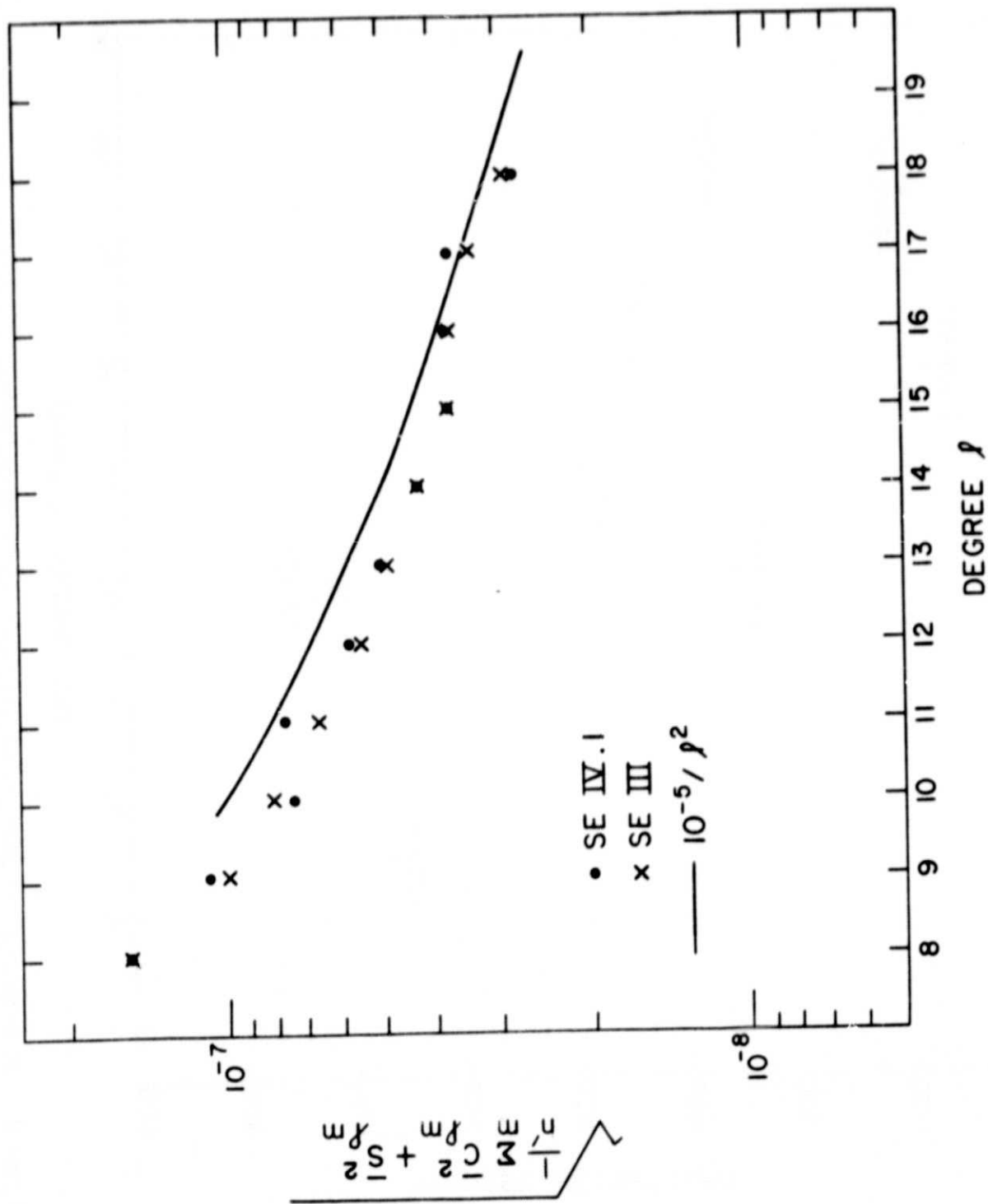


Figure 4. Mean potential coefficient by degree.

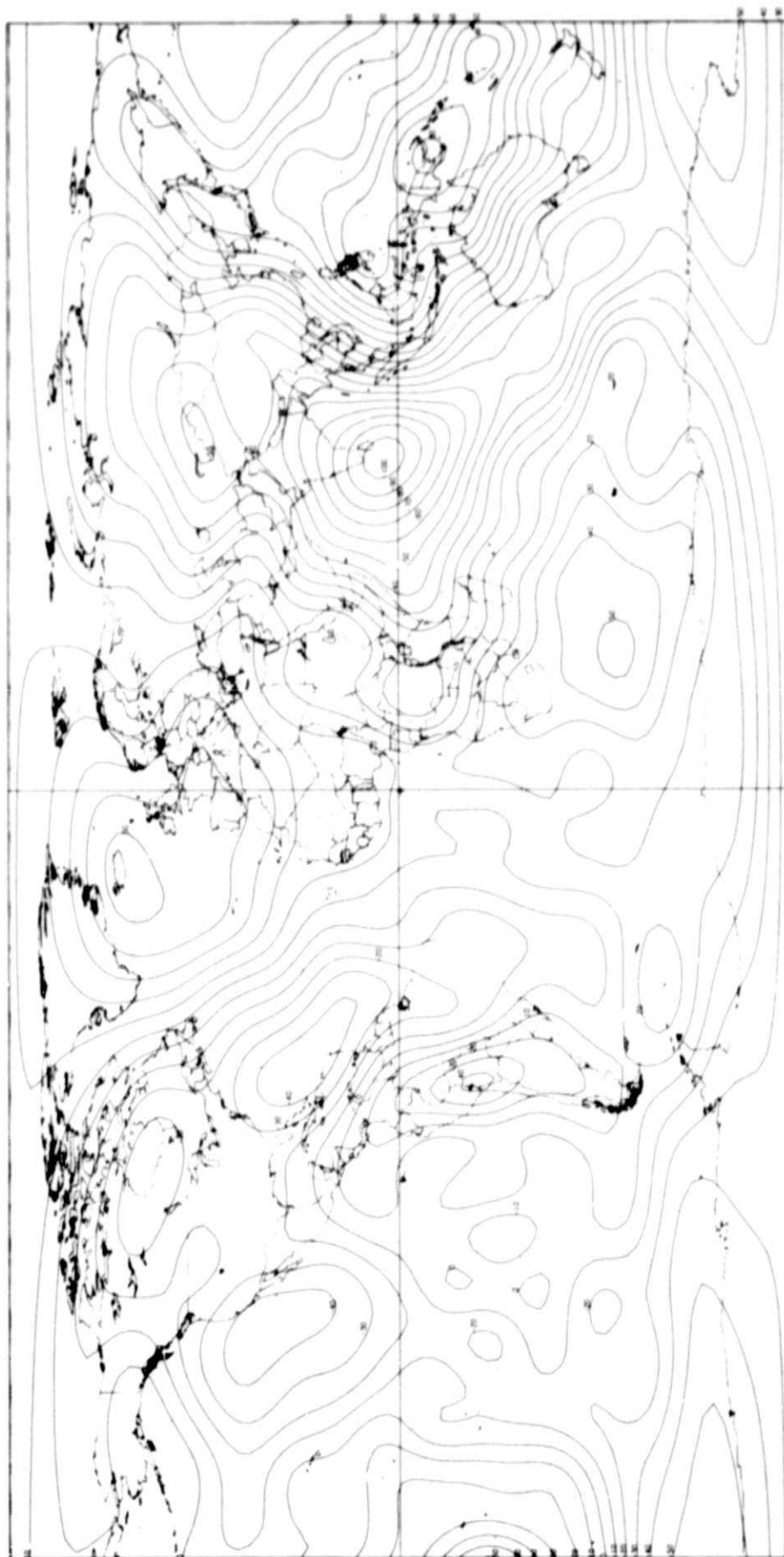


Figure 5. SE IV.1 geoid height in meters calculated with respect to the best-fitting ellipsoid, $f = 1/298.256$

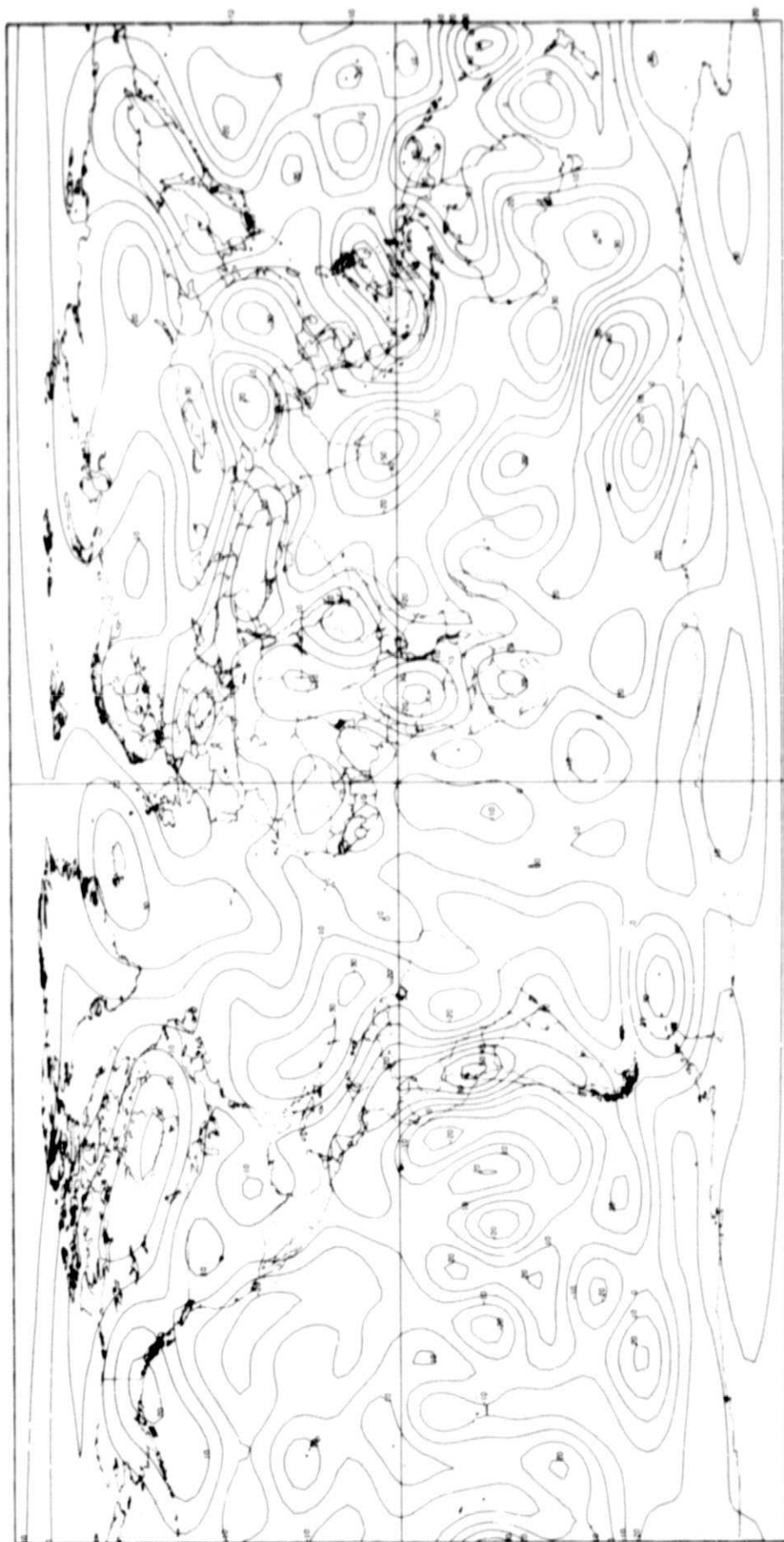


Figure 6. SE IV.1 gravity anomalies in milligals calculated with respect to the best-fitting ellipsoid, $f = 1/298.256$